# Pieces of the Puzzle: <br> Reaching QCD on Quantum Computers <br> Hank Lamm 



## Thanks to the NuQS team!

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## Outline

1 Quantum Leap

2 Digitization
3 Initialization

4 Propagation

5 Evaluation

6 Conclusions


## Finite-Density and Real-Time QFT have sign problems



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$$
\langle\mathcal{O}\rangle=\frac{\int \mathcal{D} \phi e^{-i S_{I}} \mathcal{O} e^{-S_{R}}}{\int \mathcal{D} \phi e^{-S_{R}}} \frac{\int \mathcal{D} \phi e^{-S_{R}}}{\int \mathcal{D} \phi e^{-S_{R}} e^{-i S_{I}}}=\frac{\left\langle\mathcal{O} e^{-i S_{I}}\right\rangle_{S_{R}}}{\langle\sigma\rangle_{S_{R}}}
$$

- For real $t:\langle\sigma\rangle_{S_{R}}=0$
- For $\mu \neq 0$ : Need $\propto\langle\sigma\rangle_{S_{R}}^{-2}$ configurations


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Digital QC provide entangled bits and gates, not field theories.

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- Digitize: How are (continuous) fields represented as a register?
- Initalize: How can registers be set to a field configuration?
- Propagate: How can gates be combined to evolve states?
- Evaluate: How can observables of interest be computed?


## Army you have, Army you might have, Army you want

|  | $N_{\|q\rangle}<500$ | $N_{\|q\rangle} \rightarrow \infty$ |
| :---: | :---: | :---: |
| $N_{\mathcal{U}}$ | NISQ | NESQ |
| $\lesssim 100 N_{\|q\rangle}$ | Noisy, Interm. | Noisy, Enorm. |
| $N_{\mathcal{U}}$ | FISQ | FESQ |
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[1]
Nam, Y., J.-S. Chen, N. C. Pisenti, K. Wright, C. Delaney, D. Maslov, K. R. Brown, S. Allen, J. M. Amini, J. Apisdorf, et al. In: arXiv preprint arXiv:1902.10171 (2019).

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- Moore's law like behavior "could" render methods irrelevant.


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a & b \\
c & d
\end{array}\right)
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## Fuzzy spheres can reproduce low-lying spectrum exactly ${ }^{[6]}$


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- The $O(3)$ sigma-model is defined by the Hamiltonian

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\begin{equation*}
\mathcal{H}=\sum_{r}\left[\frac{g^{2}}{2} \boldsymbol{\pi}(r)^{2}+\frac{1}{2 g^{2} \Delta x^{2}}(\mathbf{n}(r+1)-\mathbf{n}(r))^{2}\right], \tag{1}
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\psi(\mathbf{n})=\psi_{0}+\psi_{i} n_{i}+\frac{1}{2} \psi_{i j} n_{i} n_{j}+\ldots \tag{2}
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\Psi=\psi_{0} \mathbb{1}+\psi_{i} \mathbb{J}_{i}+\frac{1}{2} \psi_{i j} \mathbb{J}_{i} \mathbb{J}_{j}+\ldots \tag{3}
\end{equation*}
$$

where $\mathbb{J}_{i}, i=1,2,3$ are generators of $S U(2)$ in a given representation $j$

[^2]
## 2 qubits per site, $12 L T / \Delta t$ CNOT gates



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- For one $S U(3)$ gauge link, we could do a $\approx 5^{3}$ lattice of $\Sigma_{1080}$
- But Wilson Action freezes at $\beta_{c} \approx 3.94(2)$ on $2^{4}$ !



## Blast from the past ${ }^{[8]}$

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- Moore's Law + Bad News

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Seem to reach $\beta_{S U(3)} \approx 6$


## What are the states of strongly-coupled theories?

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What is the proton state in terms of quarks and gluons?


## $E \rho O Q$ : A hybrid quantum-classical technique

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- Signal to noise problem, Sign problem?


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- $2 \mathbb{C}$-register Inner Product gate:

$$
\left\langle\tilde{\phi}_{1} \tilde{\phi}_{2}\right| \mathfrak{U}_{\langle\cdot,\rangle}(\theta)\left|\phi_{1} \phi_{2}\right\rangle=\delta_{\phi_{1}}^{\tilde{\phi}_{1}} \delta_{\phi_{2}}^{\tilde{\phi}_{2}} e^{i \theta\left[\phi_{2}^{\dagger} \phi_{1}+\phi_{1}^{\dagger} \phi_{2}\right]}
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## Evaluation of correlators is nontrivial ${ }^{[10]}$

[^5] 113 (2 2014).

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Expectation value of a unitary operator $U$ in a given state $|\Psi\rangle$. Introducing a single ancillary qubit, we construct a controlled unitary operator $U_{C}$, defined by

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U_{C}|\Psi\rangle|0\rangle=|\Psi\rangle|0\rangle \text { and } U_{C}|\Psi\rangle|1\rangle=U|\Psi\rangle|1\rangle \tag{4}
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The operator is not unitary, so cant be evaluated by means described above. Introduce a time-dependent perturbation of the Hamiltonian:

$$
\begin{equation*}
H_{\epsilon_{1}, \epsilon_{2}}(\tau)=H_{0}+\epsilon_{2} \delta(\tau-t) W_{\mu^{\prime} \nu^{\prime}}\left(x^{\prime}\right)+\epsilon_{1} \delta(\tau) W_{\mu \nu}(x) \tag{7}
\end{equation*}
$$

[^9]
## Evaluation of correlators is nontrivial ${ }^{[10]}$

Expectation value of a unitary operator $U$ in a given state $|\Psi\rangle$. Introducing a single ancillary qubit, we construct a controlled unitary operator $U_{C}$, defined by

$$
\begin{equation*}
U_{C}|\Psi\rangle|0\rangle=|\Psi\rangle|0\rangle \text { and } U_{C}|\Psi\rangle|1\rangle=U|\Psi\rangle|1\rangle \tag{4}
\end{equation*}
$$

Generally, the expectation value of $U$ has both real and imaginary parts.

$$
\begin{equation*}
(\langle\Psi| \otimes\langle+|) U_{C}^{\dagger}\left(\mathbb{1} \otimes \sigma_{x}\right) U_{C}(|\Psi\rangle \otimes|+\rangle)=\operatorname{Re}\langle\Psi| U|\Psi\rangle \tag{5}
\end{equation*}
$$

With this procedure in mind, how to compute a correlator of the form

$$
\begin{equation*}
\langle\Psi| \mathcal{U}(-t) W_{\mu^{\prime} \nu^{\prime}}\left(x^{\prime}\right) \mathcal{U}(t) W_{\mu \nu}(x)|\Psi\rangle \tag{6}
\end{equation*}
$$

The operator is not unitary, so cant be evaluated by means described above. Introduce a time-dependent perturbation of the Hamiltonian:

$$
\begin{equation*}
H_{\epsilon_{1}, \epsilon_{2}}(\tau)=H_{0}+\epsilon_{2} \delta(\tau-t) W_{\mu^{\prime} \nu^{\prime}}\left(x^{\prime}\right)+\epsilon_{1} \delta(\tau) W_{\mu \nu}(x) \tag{7}
\end{equation*}
$$

Time evolving forward in time with $H_{\epsilon_{1}, \epsilon_{2}}$, and back with $H_{0}$ gives $C\left(\epsilon_{1}, \epsilon_{2}\right) \equiv\langle\Psi| \mathcal{U}(-t) \mathcal{U}_{\epsilon_{1}}, \epsilon_{2}(t)|\Psi\rangle$. Differentiating twice

$$
\begin{equation*}
-\left.\frac{\partial^{2} C\left(\epsilon_{1}, \epsilon_{2}\right)}{\partial \epsilon_{1} \partial \epsilon_{2}}\right|_{\epsilon_{1}=\epsilon_{2}=0}=\left\langle\mathcal{U}(-t) W_{\mu^{\prime} \nu^{\prime}}\left(x^{\prime}\right) \mathcal{U}(t) W_{\mu \nu}(x)\right\rangle \tag{8}
\end{equation*}
$$

[^10]
## Results for $2+1 \mathrm{D} D_{4}$ gauge theory



Four $D_{4}$ registers, and uses a total of 14 qubits: 12 for physical degrees of freedom, and 2 ancillary qubits. $t=10$ with a Trotterization step of $\Delta t=0.2$. In total, the quantum simulation entailed $\sim 200$ entangling gates per Trotterization time step.


## How to obtain parton distribution functions?

$$
\begin{equation*}
f(\xi)=\int_{\infty}^{\infty} \frac{d t}{2 \pi} e^{-i \xi(n \cdot P)}\langle P| \bar{\psi}\left(t n^{\mu}\right) \gamma^{+} W_{n} \psi(0)|P\rangle \tag{9}
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Simplfy to $1+1$ Thirring, then the matrix element

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\begin{equation*}
\langle P| \chi^{\dagger}\left(t n^{\mu}\right) \chi(0)|P\rangle=\langle P| e^{i H t} \chi^{\dagger}(y) e^{-i H t} \chi(0)|P\rangle=\sum_{i, j=\{x, y\}} \frac{c_{i j}}{4}\langle P| U_{i, j}|P\rangle \tag{10}
\end{equation*}
$$

in K-S prochedure $\chi \propto \sigma_{+}$and $\chi^{\dagger} \propto \sigma_{-}$which can only be measured by decomposing into $\sigma_{x}$ and $\sigma_{y}$ measurements, so need 4 simulations where

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U_{i, j}=e^{i H t} \sigma_{i} e^{-i H t} \sigma_{j} \tag{11}
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With this Hermitian construction, we can use the same $U_{C}$ based procedure prevent collapse after first measurement $\sigma_{j}$ at the cost of $2 \times$ the measurements so 8 calculations per matrix element.

## Ongoing Work of NuQS Collaboration

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- Quantum Compilers
- Remember before FORTRAN?



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- Quantum Compilers
- Remember before FORTRAN?
- Digitize Gauge Theories
- Efficent Approximations?
- Initialize w/ Lattice Field Theory
- Avoid the state specificiation?
- Evaluate Composite matrix elements

- Parton Distribution Functions?


[^0]:    [6] Alexandru, A., P. F. Bedaque, H. Lamm, and S. Lawrence. In: (2019). arXiv: 1903.06577 [hep-lat].

[^1]:    [6] Alexandru, A., P. F. Bedaque, H. Lamm, and S. Lawrence. In: (2019). arXiv: 1903.06577 [hep-lat].

[^2]:    [6] Alexandru, A., P. F. Bedaque, H. Lamm, and S. Lawrence. In: (2019). arXiv: 1903.06577 [hep-lat].

[^3]:    [8] Bhanot, G. In: Phys. Lett. 108B (1982).

[^4]:    [9] Lamm, H. and S. Lawrence. In: Phys. Rev. Lett. 121 (2018). arXiv: 1806.06649 [quant-ph].

[^5]:    [10] Pedernales, J. S., R. Di Candia, I. L. Egusquiza, J. Casanova, and E. Solano. In: Phys. Rev. Lett,

[^6]:    [10] Pedernales, J. S., R. Di Candia, I. L. Egusquiza, J. Casanova, and E. Solano. In: Phys. Rev. Lett. 113 (2 2014).

[^7]:    [10] Pedernales, J. S., R. Di Candia, I. L. Egusquiza, J. Casanova, and E. Solano. In: Phys. Rev. Lett. 113 (2 2014).

[^8]:    [10] Pedernales, J. S., R. Di Candia, I. L. Egusquiza, J. Casanova, and E. Solano. In: Phys. Rev. Lett. 113 (2 2014).

[^9]:    [10] Pedernales, J. S., R. Di Candia, I. L. Egusquiza, J. Casanova, and E. Solano. In: Phys. Rev. Lett. 113 (2 2014).

[^10]:    [10]
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