Pieces of the Puzzle: Reaching QCD on Quantum Computers Hank Lamm



Pieces of the Puzzle

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THE GEORGE WASHINGTON UNIVERSITY

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Outline

1 Quantum Leap

2 Digitization

- 3 Initialization
- **4** Propagation
- 5 Evaluation
- 6 Conclusions



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$$\langle \mathcal{O}
angle = rac{\int \mathcal{D}\phi \ e^{-iS_I} \mathcal{O}e^{-S_R}}{\int \mathcal{D}\phi \ e^{-S_R}} rac{\int \mathcal{D}\phi \ e^{-S_R}}{\int \mathcal{D}\phi \ e^{-S_R}e^{-iS_I}} = rac{\left\langle \mathcal{O}e^{-iS_I}
ight
angle_{S_R}}{\left\langle \sigma
ight
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• For real t: $\langle \sigma \rangle_{S_B} = 0$

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For real t: ⟨σ⟩_{S_R} = 0
For μ ≠ 0: Need ∝ ⟨σ⟩⁻²_{S_R} configurations

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Converting Bits to Qubits



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Digital QC provide entangled bits and gates, not field theories.

Lots of \$\$, Lots of Interest, Lots of Hype

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• Propagate: How can gates be combined to evolve states?

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$$|q\rangle^{N} \to \phi? \qquad \mathcal{U}|q\rangle^{N} \to |\psi_{0}\rangle?$$

$$\mathcal{U}|\psi_{0}\rangle \to |\psi(t)\rangle? \qquad \langle\mathcal{O}\rangle?$$

- Digitize: How are (continuous) fields represented as a register?
- Initialize: How can registers be set to a field configuration?
- Propagate: How can gates be combined to evolve states?
- Evaluate: How can observables of interest be computed?

	$N_{ q\rangle} < 500$	$N_{ q\rangle} \to \infty$
$N_{\mathcal{U}}$	NISQ	NESQ
$\lesssim 100 N_{ q\rangle}$	Noisy, Interm.	Noisy, Enorm.
$N_{\mathcal{U}}$	FISQ	FESQ
$ ightarrow\infty$	Faithful, Interm.	Faithful, Enorm.

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- Be wary of how optimizations for one era hamstring in others
- Moore's law like behavior "could" render methods irrelevant.

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• Fermions are "trival" - Bosonic fields require thought

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- Fund. Rep. with floats (Yell at Ciaran Hughes^[2])

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

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- Discrete Subgroups (Accost Me)



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Fuzzy spheres can reproduce low-lying spectrum $exactly^{[6]}$



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• Truncate the communting algebra of functions by a non-communting algebra



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Fuzzy spheres can reproduce low-lying spectrum exactly^[6]

- Truncate the communting algebra of functions by a non-communting algebra
- The O(3) sigma-model is defined by the Hamiltonian



$$\mathcal{H} = \sum_{r} \left[\frac{g^2}{2} \pi(r)^2 + \frac{1}{2g^2 \Delta x^2} (\mathbf{n}(r+1) - \mathbf{n}(r))^2 \right], \quad (1)$$

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$$\psi(\mathbf{n}) = \psi_0 + \psi_i n_i + \frac{1}{2} \psi_{ij} n_i n_j + \dots \qquad (2)$$

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$$\psi(\mathbf{n}) = \psi_0 + \psi_i n_i + \frac{1}{2} \psi_{ij} n_i n_j + \dots$$
 (2)

$$\Psi = \psi_0 \mathbb{1} + \psi_i \mathbb{J}_i + \frac{1}{2} \psi_{ij} \mathbb{J}_j \mathbb{1}_j + \dots, \qquad (3)$$

where \mathbb{J}_i , i = 1, 2, 3 are generators of SU(2) in a given representation j

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2 qubits per site, $12LT/\Delta t$ CNOT gates



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SU(3) link: 9 complex-valued double-precision floats
 → 9 × 2 × 64 = 1152 bits

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- SU(3) link: 9 complex-valued double-precision floats
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- Σ_{1080} might require 11 qubits per link.
- For one SU(3) gauge link, we could do a $\approx 5^3$ lattice of Σ_{1080}
- But Wilson Action freezes at $\beta_c \approx 3.94(2)$ on 2^4 !



Blast from the past^[8]

[8] Bhanot, G. In: Phys. Lett. 108B (1982).

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Blast from the $past^{[8]}$

• Why chose the Wilson Action?

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• Why chose the Wilson Action?

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$$S = \sum \frac{\beta_F}{6} \operatorname{Tr} U + \frac{\beta_A}{9} |\operatorname{Tr} U|^2$$

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• Why chose the Wilson Action?

• $S = \sum \frac{\beta_F}{6} \operatorname{Tr} U +$ $\frac{\beta_A}{q} |\operatorname{Tr} U|^2$

•
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• Extrapolating from 1^{st} order line/ $g^2 = 1$ lines

Blast from the $past^{[8]}$



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- Extrapolating from 1^{st} order line/ $q^2 = 1$ lines
- Potts Model: $\frac{d\beta_A}{d\beta_E} \approx 1.26$
- Moore's Law + Bad News

$S = \sum \frac{\beta_0}{6} \operatorname{Tr} U + \beta_1 \operatorname{Tr} U^2$

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May 3, 2019 1

$S = \sum \frac{\beta_0}{6} \operatorname{Tr} U + \beta_1 \operatorname{Tr} U^2$

Seem to reach $\beta_{SU(3)} \approx 6$



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What are the states of strongly-coupled theories?

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What is the proton state in terms of quarks and gluons?



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Lamm, H. and S. Lawrence. In: Phys. Rev. Lett. 121 (2018). arXiv: 1806.06649 [quant-ph].

[9]

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• Combine resources to solve nonequilibrium dynamics of many-body quantum systems^[9]

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- Quantum: Time-evolve elements of ρ as pure states

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Heisenberg Spin Chain in Magnetic Fields



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 $\langle m_x(t) \rangle$ for a N = 5 with $\mu_x(0) = 1$, $\beta = 1$, and $\mu_x(t > 0) = -1$. Forest QVM are red circles and exact result is black line.



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 $\langle m_x(t) \rangle / \langle m_x(0) \rangle$ for N = 1, with $\mu_x(0) = 1$, $\mu_z(0) = 1$, $\beta = 1.0$, and $\mu_x(t > 0) = -1$. Agave are red circles and exact result is black line.

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Can we use LQCD to initialize real-time efficiently?

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Can we use LQCD to initialize real-time efficiently?

 \bullet Proposal for extending $\mathrm{E}\rho\mathrm{OQ}$ to QFT a la Schwinger-Kelydsh:

$$\langle \mathcal{O} \rangle = \frac{\mathrm{Tr} \rho_{ij} P_{jk} \mathcal{O}_{ki}}{\mathrm{Tr} \, \rho_{ij} \delta_{ji}}$$

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- Signal to noise problem, Sign problem?

What low-level primatives are required?

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- 2 C-register Inner Product gate: $\langle \tilde{\phi}_1 \tilde{\phi}_2 | \mathfrak{U}_{\langle \cdot, \cdot \rangle}(\theta) | \phi_1 \phi_2 \rangle = \delta_{\phi_1}^{\tilde{\phi}_1} \delta_{\phi_2}^{\tilde{\phi}_2} e^{i\theta \left[\phi_2^{\dagger} \phi_1 + \phi_1^{\dagger} \phi_2 \right]}$

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^[10] Pedernales, J. S., R. Di Candia, I. L. Egusquiza, J. Casanova, and E. Solano. In: Phys. Rev. Lett. 113 (2 2014).

Expectation value of a unitary operator U in a given state $|\Psi\rangle$. Introducing a single ancillary qubit, we construct a controlled unitary operator U_C , defined by

$$U_C |\Psi\rangle |0\rangle = |\Psi\rangle |0\rangle$$
 and $U_C |\Psi\rangle |1\rangle = U |\Psi\rangle |1\rangle$. (4)

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Generally, the expectation value of U has both real and imaginary parts.

$$\left(\left\langle\Psi\right|\otimes\left\langle+\right|\right)U_{C}^{\dagger}\left(\mathbb{1}\otimes\sigma_{x}\right)U_{C}\left(\left|\Psi\right\rangle\otimes\left|+\right\rangle\right) = \operatorname{Re}\left\langle\Psi\right|U\left|\Psi\right\rangle\tag{5}$$

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With this procedure in mind, how to compute a correlator of the form

$$\langle \Psi | \mathcal{U}(-t) W_{\mu'\nu'}(x') \mathcal{U}(t) W_{\mu\nu}(x) | \Psi \rangle .$$
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The operator is not unitary, so cant be evaluated by means described above. Introduce a time-dependent perturbation of the Hamiltonian:

$$H_{\epsilon_1,\epsilon_2}(\tau) = H_0 + \epsilon_2 \delta(\tau - t) W_{\mu'\nu'}(x') + \epsilon_1 \delta(\tau) W_{\mu\nu}(x)$$

$$\tag{7}$$

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Time evolving forward in time with $H_{\epsilon_1,\epsilon_2}$, and back with H_0 gives $C(\epsilon_1,\epsilon_2) \equiv \langle \Psi | \mathcal{U}(-t)\mathcal{U}_{\epsilon_1,\epsilon_2}(t) | \Psi \rangle$. Differentiating twice

$$-\left.\frac{\partial^2 C(\epsilon_1, \epsilon_2)}{\partial \epsilon_1 \partial \epsilon_2}\right|_{\epsilon_1 = \epsilon_2 = 0} = \left\langle \mathcal{U}(-t) W_{\mu'\nu'}(x') \mathcal{U}(t) W_{\mu\nu}(x) \right\rangle \tag{8}$$

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Hank Lamm

Results for $2+1D D_4$ gauge theory



Four D_4 registers, and uses a total of 14 qubits: 12 for physical degrees of freedom, and 2 ancillary qubits. t = 10 with a Trotterization step of $\Delta t = 0.2$. In total, the quantum simulation entailed ~ 200 entangling gates per Trotterization time step.



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Pieces of the Puzzle

May 3, 2019

1/23

How to obtain parton distribution functions?

$$f(\xi) = \int_{\infty}^{\infty} \frac{dt}{2\pi} e^{-i\xi(n\cdot P)} \langle P|\bar{\psi}(tn^{\mu})\gamma^{+}W_{n}\psi(0)|P\rangle$$
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Simplify to 1+1 Thirring, then the matrix element

$$\langle P|\chi^{\dagger}(tn^{\mu})\chi(0)|P\rangle = \langle P|e^{iHt}\chi^{\dagger}(y)e^{-iHt}\chi(0)|P\rangle = \sum_{i,j=\{x,y\}} \frac{c_{ij}}{4} \langle P|U_{i,j}|P\rangle$$
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in K-S prochedure $\chi \propto \sigma_+$ and $\chi^{\dagger} \propto \sigma_-$ which can only be measured by decomposing into σ_x and σ_y measurements, so need 4 simulations where

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With this Hermitian construction, we can use the same U_C based procedure prevent collapse after first measurement σ_j at the cost of $2 \times$ the measurements so 8 calculations per matrix element.

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 - Remember before FORTRAN?



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- Initialize w/ Lattice Field Theory
 - Avoid the state specificiation?
- Evaluate Composite matrix elements
 - Parton Distribution Functions?

