Walking, σ-particle vs. dilaton, complex CFT (the sextet case study)

Lattice Higgs Collaboration (LatHC)

Zoltan Fodor, Kieran Holland, JK, Daniel Nogradi, Chik Him Wong

Julius Kuti

University of California, San Diego

Lattice BSM Workshop 2019

University of Syracuse, May 2-3, 2019

*Tantalizing big questions in the sextet case study:

Origin of near-conformal walking?

How light is the emergent 0++ scalar in chiral limit?

Need to extrapolate with IR EFT like σ -model, or dilaton, or ...

Need for new strategy of EFT in massless chiral limit

EFT in complex CFT walking paradigm?

*

review of our recent results under these tantalizing questions \rightarrow

Tantalus, a king of ancient Phrygia in Greek mythology, made the mistake of gravely offending the gods. As a punishment, once dead the king was forced to stand in a pool of water, with fruit hanging just over his head. The water would recede every time the king tried to take a sip, and the fruit would lift away every time he reached to take a bite.

origin of the emergent light 0++ scalar (σ -particle or dilaton?)



- nf=10 β -function (no IRFP is found so far, contrary to earlier claims)
- nf=12 β -function controversy includes anomalous dimension
- nf=13 β -function LatHC reports nf=13 conformal IRFP (?)
- 5-loop β -function below CW: a pair of complex conjugate zeros (complex CFT help?)
- several models close to CW focus in this talk on the sextet model

Figure 1: The conformal window from [3]. The shaded regions are the

origin of the emergent light 0++ scalar, (o-particle or dilaton?) 37



3.8

5.55 10

11

9



- two zeros of 5-loop β-function collide between nf=12 and nf=13 and turn into complex conjugate pair of zeros
- sextet model closest to CW among near-conformal (except perhaps nf=12?) complex pair of CFT?

origin of the emergent light 0++ scalar (σ -particle, or dilaton, or ...?)

simplest: light Higgs-like particle of linear σ-model in PT:

 $SU(2) \otimes SU(2) \sim O(4)$ for sextet model

 $L = \frac{1}{2} \left(\partial_{\mu} \vec{\pi} \right)^{2} + \frac{1}{2} \left(\partial_{\mu} \sigma \right)^{2} - \frac{1}{2} \mu^{2} \left(\sigma^{2} + \vec{\pi}^{2} \right) + \frac{1}{4} g \left(\sigma^{2} + \vec{\pi}^{2} \right)^{2} - \varepsilon \sigma - \varepsilon$

non-linear σ -model \rightarrow dilaton EFT:



$$V(\sigma) = \frac{1}{2}m_{\sigma}^2 \cdot \sigma^2 + d_3 \left(\frac{m_{\sigma}^2}{2f_{\pi}}\right) \cdot \sigma^3 + d_4 \left(\frac{m_{\sigma}^2}{8f_{\pi}^2}\right) \cdot \sigma^4 + \dots$$

linear σ -model limit (SM): $a = b = d_3 = d_4 = 1$ dilaton EFT will require $a = b^2$, $b_3 = 0$ and special y(μ) in L

we analyze dilaton EFT first and test if non-linear σ -model limit exists without scale-dependent anomalous dimension y

triviality analysis in $m_{\sigma}/f_{\pi} < 3$ range $m_{\pi} = 0$ circa 1987-1988

→ $m_{\sigma}^2 \ge 3m_{\pi}^2$ tree level relation $m_{\sigma}^2 \ge 2m_{\pi}^2$ with loop corrections

p-regime data: 0++ is tracking the Goldstone pion with $m_{\pi}^2 \ge m_{\sigma}^2$, not like linear σ -model New EFT is needed to extrapolate data to massless chiral limit



dilaton EFT with $\sigma(x)$ dilaton field and $\pi^{a}(x)$ Goldstone bosons

$$L = \frac{1}{2}\partial_{\mu}\chi\partial_{\mu}\chi - V_{d}(\chi) + \frac{f_{\pi}^{2}}{4}\left(\frac{\chi}{f_{d}}\right)^{2} tr\left[\partial_{\mu}\Sigma^{\dagger}\partial_{\mu}\Sigma\right] - \frac{f_{\pi}^{2}m_{\pi}^{2}}{4}\left(\frac{\chi}{f_{d}}\right)^{y} tr\left(\Sigma + \Sigma^{\dagger}\right)$$

 $y = 3 - \gamma$ where γ is the mass anomalous dimension $\chi(x) = f_d e^{\sigma(x)/f_d}$ describes the dilaton field $\sigma(x)$ pion field $\Sigma = e^{i\pi^a \tau^a/f_{\pi}}$ with τ^a Pauli matrices, tree level pion mass $m_{\pi}^2 = 2Bm$



- dilaton EFT has long history
- Golterman-Shamir expansion in x=N_f/N variable
- Veneziano limit N → ∞
- predicts walking around * in p-regime (tree level)
 from expanding around CFT *
- based on scheme-dependent β -function ?
- flavor symmetry at * is different from symmetry at * ?

dilaton EFT with $\sigma(x)$ dilaton field and $\pi^{a}(x)$ Goldstone bosons

$$L = \frac{1}{2}\partial_{\mu}\chi\partial_{\mu}\chi - V_{d}(\chi) + \frac{f_{\pi}^{2}}{4}\left(\frac{\chi}{f_{d}}\right)^{2}tr\left[\partial_{\mu}\Sigma^{\dagger}\partial_{\mu}\Sigma\right] - \frac{f_{\pi}^{2}m_{\pi}^{2}}{4}\left(\frac{\chi}{f_{d}}\right)^{y}tr\left(\Sigma + \Sigma^{\dagger}\right)$$

 $y = 3 - \gamma$ where γ is the mass anomalous dimension $\chi(x) = f_d e^{\sigma(x)/f_d}$ describes the dilaton field $\sigma(x)$ pion field $\Sigma = e^{i\pi^a \tau^a/f_{\pi}}$ with τ^a Pauli matrices, tree level pion mass $m_{\pi}^2 = 2Bm$

Golterman-Shamir

Appelquist et al. notation

but we do our own IML analysis which is required for any conclusion!

$$V_{\sigma} = \frac{m_d^2}{2f_d^2} \left(\frac{\chi^2}{2} - \frac{f_d^2}{2}\right)^2 \text{ relevant deformation of IRFP theory} \begin{cases} f_{\pi} & \text{Goldstone decay constant} \\ m_{\pi} = 2mB & \text{Goldstone pions} \\ f_d & \text{dilaton decay constant} \\ m_d & \text{dilaton mass} \end{cases}$$

$$V_{d} = \frac{m_d^2}{16f_d^2} \chi^4 \left(4\ln\frac{\chi}{f_d} - 1\right) \text{ nearly marginal deformation} \\ \text{Golterman-Shamir form} \end{cases}$$

 $M_{\pi}^{2} \cdot F_{\pi}^{2-y} - 2B_{\pi} \cdot f_{\pi}^{(2-y)} \cdot m = 0 \quad \text{general V indep. scaling law}$ $F_{\pi}^{(4-y)} \cdot (1 - f_{\pi}^{2}/F_{\pi}^{2}) - 2y \cdot n_{f} f_{\pi}^{(6-y)} B_{\pi}/m_{d}^{2} f_{d}^{2} \cdot m = 0 \quad V_{\sigma}'(\chi = F_{d})$ $3F_{\pi}^{2}/M_{\pi}^{2} - f_{\pi}^{2}/M_{\pi}^{2} - 2M_{d}^{2}/m_{d}^{2} \cdot f_{\pi}^{2}/M_{\pi}^{2} - y(y - 1)n_{f} f_{\pi}^{4}/m_{d}^{2} f_{d}^{2} = 0 \quad V_{\sigma}''(\chi = M_{d})$ $F_{\pi}^{(4-y)} \cdot \log(F_{\pi}/f_{\pi}) - y \cdot n_{f} f_{\pi}^{(6-y)} B_{\pi} \cdot m/m_{d}^{2} f_{d}^{2} = 0 \quad V_{d}'(\chi = F_{d})$ $(F_{\pi}^{2}/M_{\pi}^{2}) \cdot (3\log(F_{\pi}/f_{\pi}) + 1) - (M_{d}^{2}/m_{d}^{2}) \cdot (f_{\pi}^{2}/M_{\pi}^{2}) - y(y - 1)n_{f} f_{\pi}^{4}/2m_{d}^{2} f_{d}^{2} = 0 \quad V_{d}''(\chi = M_{d})$

 M_{π}, F_{π}, M_{d} input data at each m $f_{\pi}, B, f_{d}, m_{d}, y$ fitted for all m IML: Implicit Maximum Likelihood test IML is very different from ML fitting Perfect fits for V_{σ} ! V_{d} fails!

dilaton EFT with $\sigma(x)$ dilaton field and $\pi^{a}(x)$ Goldstone bosons (sextet model)



dilaton EFT with $\sigma(x)$ dilaton field and $\pi^{a}(x)$ Goldstone bosons (sextet model)



dilaton EFT with $\sigma(x)$ dilaton field and $\pi^{a}(x)$ Goldstone bosons (sextet model)



to reach the chiral regime requires two orders of magnitude drop in fermion mass: switch from p-regime to epsilon regime and related RMT

epsilon regime and RMT

$$\mathscr{L}_{\varepsilon} = \frac{1}{2} \partial_{\mu} \chi \partial_{\mu} \chi - V_d(\chi) + \frac{m_{\pi}^2 f_{\pi}^2}{4} \left(\frac{\chi}{f_d}\right)^{y} \operatorname{tr}\left[\Sigma_0 + \Sigma_0^{\dagger}\right]$$

 $\mathscr{L}_{\delta} = \frac{1}{2} \partial_{\mu} \chi \partial_{\mu} \chi - V(\chi) + \frac{f_{\pi}^{2}}{4} \left(\frac{\chi}{f_{d}}\right)^{2} \operatorname{tr}\left[\partial_{t} \Sigma_{0} \ \partial_{t} \Sigma_{0}^{\dagger}\right]$

epsilon regime with very small fermion mass deformation

delta regime m=0 very small fermion mass deformation can be added

new ensembles at equivalent p-regime pion mass Mpi ~ 100 and volume size 64⁴



epsilon regime and RMT



successful testing

ongoing analysis (preliminary results not shown)



Luca Vecchi 2010 talks about complex CFT built on Gies et al., Terao et al., Kaplan et al....

walking and complex CFT new paradigm? flavor symmetry group is the same for walker and the CFT! paradigm change







Luca Vecchi 2010 talks about complex CFT built on Gies et al., Terao et al., Kaplan et al...

walking and complex CFT new paradigm? flavor symmetry group is the same for walker and the CFT! paradigm change

we started work earlier on the realization of walking based on this idea

to distinguish near-conformal and conformal finite volume correlators (drifting scaling exponents distinguished from fixed conformal exponents). we ran into difficulty not knowing the conformal exponents of the complex theory.

Gorbenko, Rychkov, Zan turned to a two-dimensional example (Potts model) for very detailed realization of working without apparently knowing about Vecchi

We ask now: can we realize the lattice Potts conformal field theories for continuous Q (flavor)



Potts model 7 Q potts spin¹⁰ flavor described by CFT

$$\begin{split} \varepsilon = . & Ocrean 4 gair of real CFT with pair of zeros of the beta function works for continuous Q in cluster rep \\ & Solve \left[\left(\frac{-2 I \varepsilon}{4 \pi complex} - \frac{\varepsilon^2}{4} \right) + \pi \left(\frac{4}{4 \pi complex} - \frac{2 I \sqrt{3}}{6} \varepsilon \right) g + \frac{8 \pi^2}{3} g^2 = 0, g \right] \\ & \left\{ \left\{ g + \frac{1 \sqrt{3} \varepsilon}{4 \pi complex} + \frac{-2 \sqrt{3} \pi + i \sqrt{3} \varepsilon}{6 \pi c (1 + 1) \sqrt{3} \varepsilon} \right\} \right\} \\ & \left\{ \left\{ g + \frac{1 \sqrt{3} \varepsilon}{4 \pi complex} + \frac{-2 \sqrt{3} \pi + i \sqrt{3} \varepsilon}{6 \pi c (1 + 1) \sqrt{3} \varepsilon} \right\} \right\} \\ & \left\{ \left\{ g + \frac{1 \sqrt{3} \varepsilon}{4 \pi complex} + \frac{-2 \sqrt{3} \pi + i \sqrt{3} \varepsilon}{6 \pi c (1 + 1) \sqrt{3} \varepsilon} \right\} \\ & \left\{ \left\{ g + \frac{1 \sqrt{3} \varepsilon}{4 \pi c (1 + 1) \sqrt{3} \varepsilon} + \frac{-2 \sqrt{3} \pi + i \sqrt{3} \varepsilon}{6 \pi c (1 + 1) \sqrt{3} \varepsilon} \right\} \\ & \left\{ \left\{ g + \frac{1 \sqrt{3} \varepsilon}{4 \pi c (1 + 1) \sqrt{3} \varepsilon} + \frac{-2 \sqrt{3} \pi + i \sqrt{3} \varepsilon}{6 \pi c (1 + 1) \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{6 \pi c (1 + 1) \sqrt{3} \varepsilon} \right\} \\ & \left\{ \left\{ g + \frac{1 \sqrt{3} \varepsilon}{4 \pi c (1 + 1) \sqrt{3} \varepsilon} + \frac{-2 \sqrt{3} \pi + i \sqrt{3} \varepsilon}{6 \pi c (1 + 1) \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 + 1 \sqrt{3} \varepsilon} + \frac{1 \sqrt{3} \varepsilon}{1 +$$







RG flow of g between two complex fixed points Q=4.01 0.02 0.015 0.01 0.005 lm(g) -0.005 -0.01 -0.015 -10 -2 0 2 6 -8 -6 8 $imes 10^{-3}$ Real(g)

Potts model Q potts spin ~ flavor described by CFT

Q=2 - 4 pair of real CFT with pair of zeros of the beta function works for continuous Q in cluster rep

Q > 4 complex CFT, like Q=5, 6, 7 ...

What was identified before as Q=5 Potts is near-conformal and walking, supported by the complex IRFP pair, not critical in very well defined sense

We use efficient Potts CUDA code for integer Q Swendsen-Wang ~1 ns/spin

CPU cluster code for arbitrary Q

Tensor method and DMRG on Potts transfer matrix

working on lattice realization of the Potts complex CFT paradigm



Thank you