

Constraint Effective Potentials in Gauge-Higgs Unification models

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Motivation

Gauge-Higgs Unification

- ▶ origin of the potential responsible for the Brout-Englert-Higgs mechanism still unknown
- ▶ in Gauge-Higgs Unification (GHU) models the Higgs field is associated with some extra-dimensional components of a gauge field [Manton, 1979]
- ▶ one extra dimension: Higgs potential is zero at tree level and generated only through quantum effects
- ▶ 1-loop: Higgs mass is finite, fermions needed for SB through Hosotani mechanism [Hosotani, 1983].
- ▶ what happens non-perturbatively? \longrightarrow lattice study.
- ▶ measure constraint effective potential [Kuti, Shen, 1988]
- ▶ also relevant for effective Polyakov loop actions



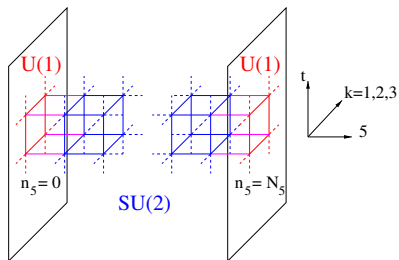
5D Orbifold Model Lattice Formulation

Wilson Gauge action for bulk gauge group $SU(2)$:

$$S_W^{orb} = \frac{\beta_4}{2} \sum_{\mu, \nu} w \cdot \text{tr} \{1 - U_{\mu\nu}\} + \frac{\beta_5}{2} \sum_{\mu} \text{tr} \{1 - U_{\mu 5}\}$$

Boundary links satisfy $gUg^{-1} = U$ with $g = -i\sigma^3$

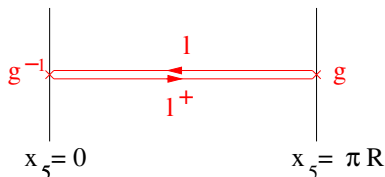
$$w = \begin{cases} \frac{1}{2} & \text{plaquette on boundary} \\ 1 & \text{otherwise} \end{cases}$$



5D Orbifold

- ▶ bare anisotropy is $\gamma = \sqrt{\beta_5/\beta_4} \simeq a_4/a_5$
- ▶ N_5 is number of *links* in the fifth dimension
- ▶ $N_5 a_5 = \pi R$, R radius of extra dimension

5D Orbifold Model Higgs operators



Scalar Polyakov loop (defined at $n = (n_\mu, 0)$)

$$p(n) = l(n) g l^\dagger(n) g^\dagger$$

Higgs field

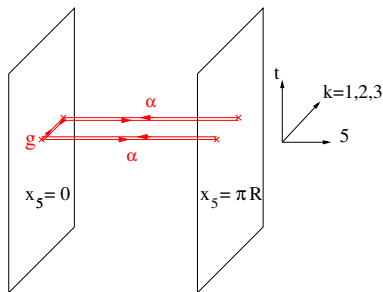
$$h(n) = [p(n) - p^\dagger(n), g]/(4N_5) \sim A_5^1 \sigma^1 + A_5^2 \sigma^2$$

Higgs operators

$$\mathcal{H}(n_0) = \sum_{n_1, n_2, n_3} \text{tr} [hh^\dagger] \quad , \quad \mathcal{P}(n_0) = \sum_{n_1, n_2, n_3} \text{tr} [p]$$



Lattice operators: Gauge boson



Gauge boson operators (defined at $n = (n_\mu, 0)$)

$$\mathcal{Z}(n_0, k) = \sum_{n_1, n_2, n_3} \text{tr} [g U(n, k) \alpha(n + a_4 \hat{k}) U^\dagger(n, k) \alpha(n)]$$

$$\mathcal{Z}'(n_0, k) = \sum_{n_1, n_2, n_3} \text{tr} [g U(n, k) l(n + a_4 \hat{k}) U^\dagger(n', k) l^\dagger(n)]$$

α is the $SU(2)$ projection of h ; $n' = (n_\mu, N_5)$



Spontaneous symmetry breaking

Elitzur's theorem [Elitzur, 1975]

In the gauge-invariant lattice theory, spontaneous symmetry breaking (SSB) must originate from the breaking of a *global* symmetry.

Stick symmetry

Global transformation [Ishiyama, Murata, So and Takenaga, 0911.4555]

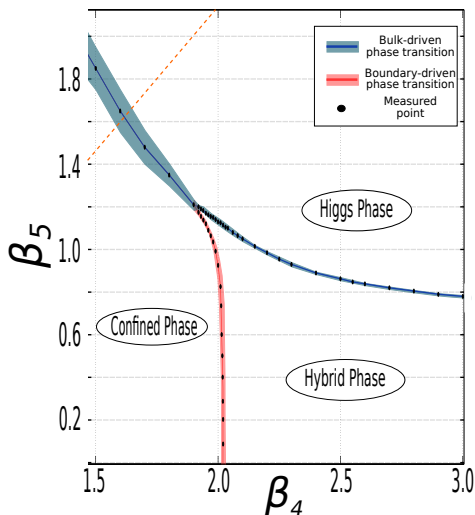
$$\begin{cases} U(n_5 = 0, 5) \rightarrow g_s^{-1} U(n_5 = 0, 5) \\ U(n_5 = 0, \mu) \rightarrow g_s^{-1} U(n_5 = 0, \mu) g_s \end{cases} \quad \text{with } g_s = -i\sigma^1$$

Gauge boson operators \mathcal{Z} , \mathcal{Z}' are odd under the stick symmetry, they are the order parameters of SSB

[Irges and Knechtli, arXiv:1312.3142]



5D Orbifold Model Phase Diagram



Notes

On the torus

- ▶ confined and Coulomb phase

On the orbifold

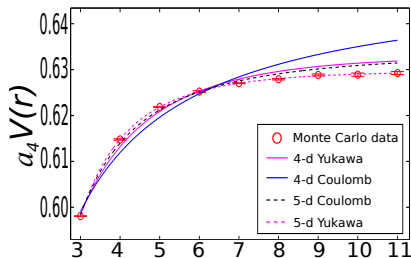
- ▶ One more phase: $U(1)$ gauge links deconfine separately
- ▶ No compactification observed at $\gamma > 1$
- ▶ Interesting physics is found at $\gamma < 1$

first order PTs [Alberti, Irges, Knechtli and Moir, 1506.06035]

Dimensional Reduction

Higgs phase

- ▶ The orbifold's $SU(2)$ bulk feels the fifth dimension

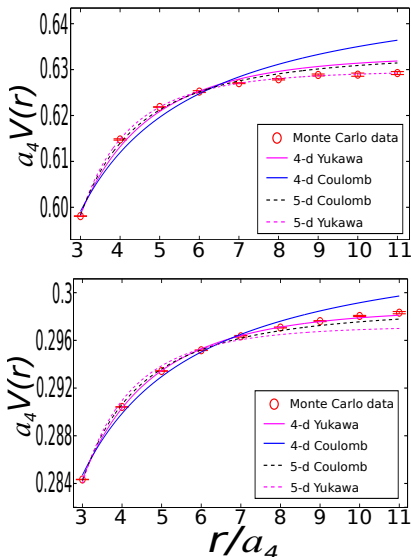


Dimensional Reduction

Higgs phase

- ▶ The orbifold's $SU(2)$ bulk feels the fifth dimension
- ▶ But near the PT the boundary remains four-dimensional
- ▶ The fitted Yukawa masses agree with the measured one

⇒ dimensional reduction
via **localization**



[Alberti, Irges, Knechtli and Moir, 1506.06035]



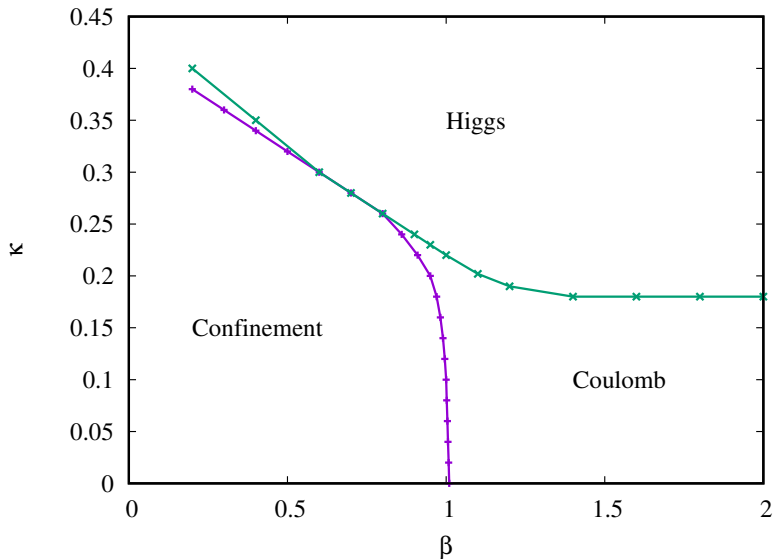
5D Orbifold Model Summary

Non-perturbative Gauge-Higgs Unification

- ▶ SU(2) pure gauge theory on a 5D orbifold has a Higgs phase with the Higgs mechanism realized as a quantum and bosonic effect
- ▶ a Standard-Model like spectrum can be reproduced
- ▶ cut-off effects appear to be small, excited state energies are 2–5 times larger than the ground states
- ▶ localization on the 4D boundaries is observed
- ▶ 5D Orbifold boundary resembles a 4D Abelian gauge-Higgs model (with $q = 2$)
- ▶
$$S_\phi[U, \phi] = \sum_x |\phi(x)|^2 + \lambda(|\phi(x)|^2 - 1)^2 - 2\kappa \sum_\mu \text{Re} \left\{ \phi^\dagger(x) [U_\mu(x)]^q \phi(x + a\hat{\mu}) \right\}$$
- ▶ first study of U_Ω in a gauge theory...



4D Abelian Higgs Phase Diagram ($\lambda = 1.0$)



4D Abelian Higgs Unitary Gauge

$$\phi(x) = \rho(x) \exp i\varphi(x) \rightarrow \phi_1 = \rho \cos \varphi, \phi_2 = \rho \sin \varphi$$

$$\Rightarrow d\phi_1 d\phi_2 = \left| \det \begin{pmatrix} \cos \varphi & -\rho \sin \varphi \\ \sin \varphi & \rho \cos \varphi \end{pmatrix} \right| d\rho d\varphi = \rho d\rho d\varphi$$

$$S_\rho[V, \rho] = \sum_x \rho_x^2 + \lambda(\rho_x^2 - 1)^2$$

$$-2\kappa\rho_x \sum_\mu \rho_{x+\hat{\mu}} \underbrace{\text{tr} (e^{-i\varphi_x} U_{x,\mu} e^{i\varphi_{x+\hat{\mu}}})}_{=V_{x,\mu}}$$

$$\prod_{x,\mu} \int_{SU(2)} dU_{x,\mu} \prod_x \int_{-\infty}^{\infty} d\phi(x) e^{-S(U,\phi)} \rightarrow$$

$$\prod_{x,\mu} \int_{SU(2)} dU_{x,\mu} \prod_x \int_0^\infty \rho(x) d\rho(x) \int_{U(1)} d\varphi(x) e^{-S(U,\rho,\varphi)} \propto$$

$$\prod_{x,\mu} \int_{SU(2)} dV_{x,\mu} \prod_x \int_0^\infty d\rho(x) e^{-S(V,\rho) + \ln \rho}$$



4D Abelian Higgs Constraint HMC

$$H[V, \rho] = S_\rho[V, \rho] - \ln \rho + \frac{1}{2} \sum_x \pi(x)^2 + \mu \left(\frac{1}{\Omega} \sum_x \rho(x) - \Phi \right)$$

constraint EOMs: [[Fodor, Holland, Kuti, Nogradi, Schroeder, 0710.3151](#)]

$$\dot{\rho}(x, t) = \frac{\partial H}{\partial \pi(x, t)} = \pi(x, t)$$

$$\dot{\pi}(x, t) = -\frac{\partial H}{\partial \rho(x, t)} = -\frac{\partial S_\rho}{\partial \rho(x, t)} + \frac{1}{\rho} - \frac{\mu}{\Omega}$$

time derivatives of the constraint to ensure it's unchanged

$$\frac{\partial}{\partial t} \left(\frac{1}{\Omega} \sum_x \rho(x) - \Phi \right) = \sum_x \dot{\rho}(x) = \sum_x \pi(x) = 0$$

$$\frac{\partial}{\partial t} \sum_x \pi(x) = \sum_x \dot{\pi}(x) = 0 \Rightarrow \mu = \sum_x \left(\frac{1}{\rho} - \frac{\partial S_\rho}{\partial \rho(x)} \right)$$



4D Abelian Higgs Constraint Effective Potential

the derivative of the potential with respect to Φ can be obtained from

$$\begin{aligned}
 e^{-\Omega U_{\Omega}(\Phi)} &= \int \mathcal{D}\rho \mathcal{D}V \delta\left(\frac{1}{\Omega} \sum_x \rho(x) - \Phi\right) e^{-S + \ln(\rho)} \\
 &\propto - \int \mathcal{D}\rho \mathcal{D}V \mathcal{D}\pi e^{-H[\rho, V, \Phi]}
 \end{aligned}$$

using a variable shift $\rho = \rho' + \Phi$ to avoid the δ -function in the first line [O'Raifeartaigh, Wipf, Yoneyama, 1986], or equivalently

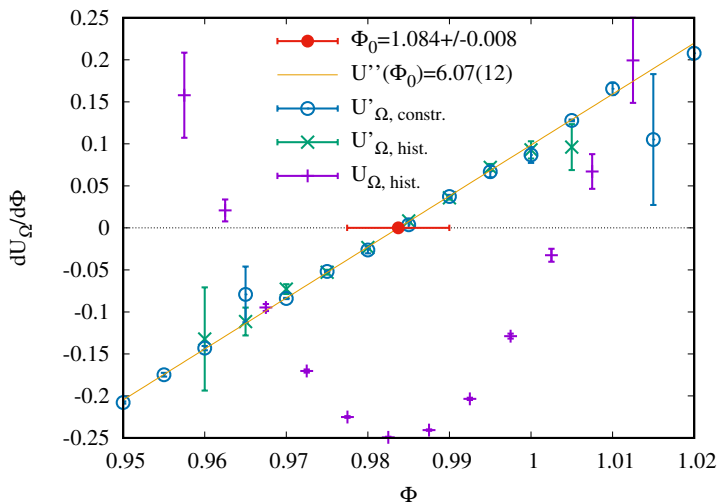
$$\begin{aligned}
 -\Omega U'_{\Omega}(\Phi) e^{-\Omega U_{\Omega}(\Phi)} &= - \int \mathcal{D}\rho \mathcal{D}V \mathcal{D}\pi e^{-H[\rho, V, \Phi]} \frac{dH}{d\Phi} \\
 \Rightarrow U'_{\Omega}(\Phi) &= \frac{1}{\Omega} \left\langle \frac{dH}{d\Phi} = -\mu \right\rangle_{\Phi} = \frac{1}{\Omega} \left\langle \sum_x \left(\frac{\partial S_{\rho}}{\partial \rho(x)} - \frac{1}{\rho} \right) \right\rangle_{\Phi}
 \end{aligned}$$

where $\langle \dots \rangle_{\Phi}$ means the expectation value at fixed Φ



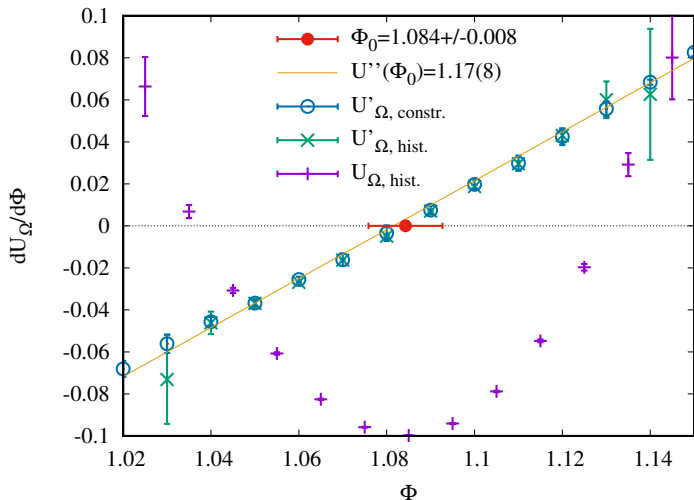
4D Abelian Higgs Effective Potential

$$\beta = 0.6, \kappa = 0.3, \lambda = 1.0, am_H \approx 2.1(3)$$



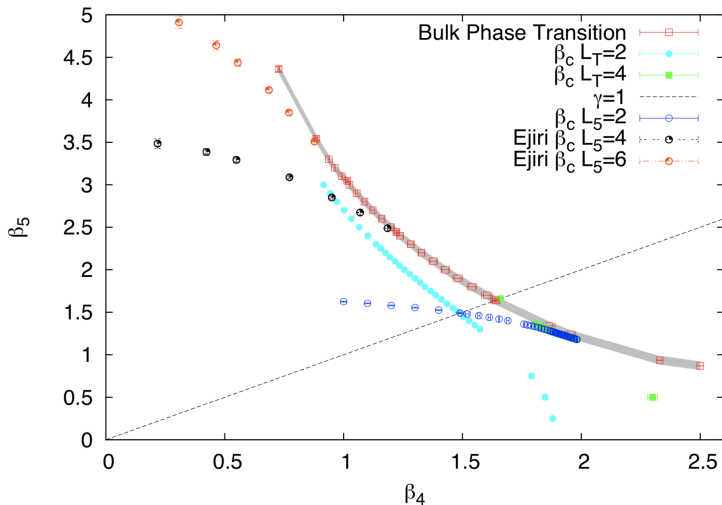
4D Abelian Higgs Effective Potential

$$\beta = 1.4, \kappa = 0.17, \lambda = 0.15, am_H \approx 0.9(2)$$



5D Torus Phase Diagram [Knechtli, Luz, Rago, arXiv:1110.4210]

first step towards 5D orbifold model we work on the torus



5D Torus Constraint HMC

new gauge-invariant, symplectic algorithm for constraint Hamiltonian of the Higgs operator given by the Polyakov loop along the fifth dimension $P_5 = \prod_{n=0}^{N_5} [U_5(x, n)] = e^{N_5 H}$

$$S_W^{tor} = \frac{\beta_4}{2} \sum_{\mu, \nu} \text{tr} \{1 - U_{\mu\nu}\} + \frac{\beta_5}{2} \sum_{\mu} \text{tr} \{1 - U_{\mu 5}\}$$

$$H[U] = S_W^{tor}[U] + \frac{1}{2} \sum_{\mathbf{x}, \mu} \text{tr} [\pi_{\mu}^2(\mathbf{x})] + \lambda \left(\frac{1}{2\Omega} \sum_x \text{tr} P_5(x) - \Phi \right)$$

$$\dot{U}_5(x, n_5) = \pi_5(x, n_5) U_5(x, n_5)$$

$$\dot{\pi}_5(x, n_5) = -\frac{\partial S[U_5]}{\partial U_5(x, n_5)} + \frac{\lambda}{8\Omega} \text{tr} [\dots \sigma_i U_5(x, n_5) \dots] \sigma^i$$

following [Hairer, Lubich, Wanner, 2006] we use an extension of the so-called Rattle algorithm to general Hamiltonians for constraint systems and apply it to our problem



Newton-Störmer-Verlet-leapfrog method on the Torus

$$\pi_{n+1/2} = \pi_n - \frac{h}{2} \left(\frac{\partial S}{\partial U_n} - \frac{\lambda}{8\Omega} \text{tr} [\dots \sigma_i U_n \dots] \sigma^i \right)$$

$$U_{n+1} = e^{h\pi_{n+1/2}} U_n = [1 + h\pi_n - h^2/2(\dots + \pi_n^2) + \mathcal{O}(h^3)] U_n$$

$$0 = \frac{1}{2\Omega} \sum_x \text{tr} P_{n+1} - \Phi = \frac{1}{2\Omega} \sum_x \text{tr} \prod_{n_5=0}^{N_5-1} U_{n+1}(x, n_5) - \Phi$$

$$\pi_{n+1} = \pi_{n+1/2} - \frac{h}{2} \left(\frac{\partial S}{\partial U_{n+1}} - \frac{\mu}{8\Omega} \text{tr} [\dots \sigma_i U_{n+1} \dots] \sigma^i \right)$$

$$0 = \frac{1}{8\Omega} \sum_{x, n_5} \text{tr} \{ \text{tr} [\dots \sigma_i U_{n+1}(x, n_5) \dots] \sigma^i \pi_{n+1}(x, n_5) \}$$

the first three equations determine $(\pi_{n+1/2}, U_{n+1}, \lambda)$,
 whereas the remaining two give (π_{n+1}, μ)

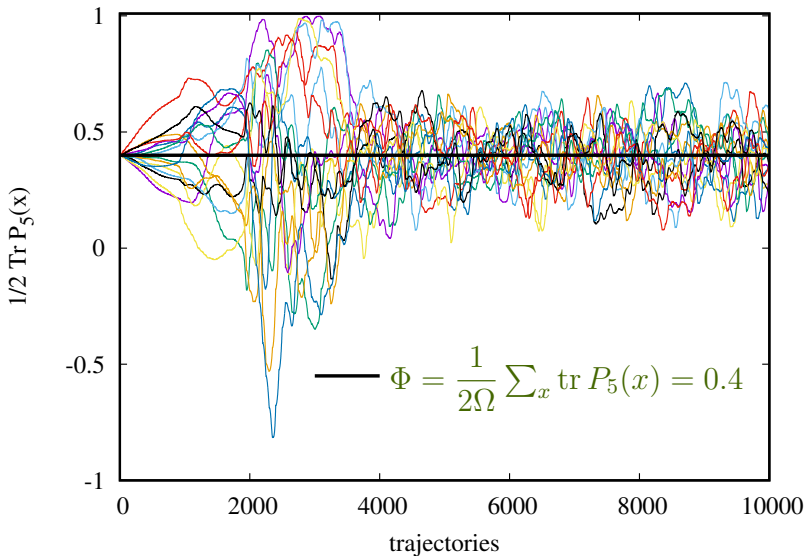


Newton-Störmer-Verlet-leapfrog method on the Torus

$$\begin{aligned}
 \frac{\lambda}{8\Omega} &= \left\{ \sum_{x, n_5} \left(\text{tr} \left[\dots \left(\frac{\partial S}{\partial U_n(x, n_5)} - 2\pi_n(x, n_5)/h \right) U_n(x, n_5) \dots \right] \right. \right. \\
 &\quad - 2 \sum_{m_5 > n_5}^{N_5-1} \text{tr} \left[\dots \pi_n(x, n_5) U_n(x, n_5) \dots \pi_n(x, m_5) U_n(x, m_5) \dots \right] \\
 &\quad \left. \left. - \text{tr} \left[\dots \pi_n^2(x, n_5) U_n(x, n_5) \dots \right] \right) \right\} / \\
 &\quad \sum_{x, n_5} \text{tr} \left\{ \dots \text{tr} \left[\dots \sigma_i U_n(x, n_5) \dots \right] \sigma^i U_n(x, n_5) \dots \right\} + \mathcal{O}(h^3)!!! \\
 \frac{\mu}{8\Omega} &= \sum_{x, n_5} \left(\text{tr} \left[\dots \sigma_i U_{n+1}(x, n_5) \dots \right] \text{tr} \left[\sigma^i \partial S / \partial U_{n+1}(x, n_5) \right] \right. \\
 &\quad \left. - 2 \text{tr} \left[\dots \sigma_i U_{n+1}(x, n_5) \dots \right] \text{tr} \left[\sigma^i \pi_{n+1/2}(x, n_5) \right] / h \right) / \\
 &\quad \sum \text{tr} \left\{ \left(\text{tr} \left[\dots \sigma_i U_{n+1}(x, n_5) \dots \right] \sigma^i \right)^2 \right\}
 \end{aligned}$$

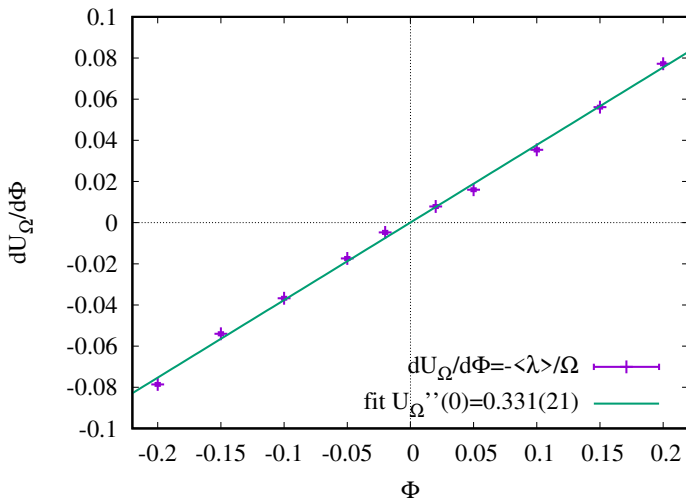


Newton-Störmer-Verlet-leapfrog method on the Torus



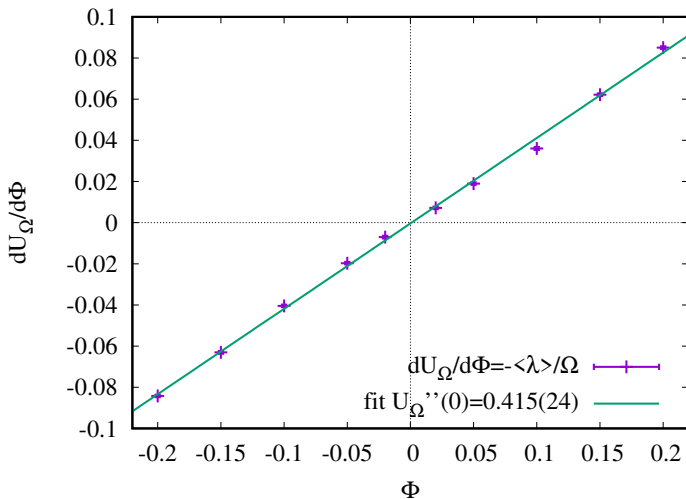
5D Torus Effective Higgs Potential

$$\beta_4 = 1.0, \beta_5 = 2.8, \Omega = 8^4, N_5 = 4, \langle \text{tr } P_5 \rangle = 0, a_4 m_H \approx 0.5(1)$$



5D Torus Effective Higgs Potential

$$\beta_4 = 0.8, \beta_5 = 3.8, \Omega = 8^4, N_5 = 6, \langle \text{tr } P_5 \rangle = 0, a_4 m_H \approx 0.7(1)$$



Conclusions & Outlook

Constraint HMC Algorithms for gauge-Higgs models

- ▶ implemented the constraint HMC for 4D Abelian gauge-Higgs model and computed the effective Higgs potential in the spontaneously broken phase
- ▶ developed a new symplectic constraint HMC algorithm for 5D torus and found no SSB in accordance with the absence of stick symmetry

Open Questions and Outlook

- ▶ interpretation of torus results, $P_5 = e^{N_5 H}$
- ▶ constraint HMC algorithm for 5D orbifold $\Phi = \langle \text{tr } \mathcal{Z} \rangle$
- ▶ other applications: (constraint) effective Polyakov loop action for finite temperature QCD

