

## 7320 FINAL EXAM

Begin each problem on a separate piece of paper. Take your time and think before you write. Show all your work. When an explanation is required, write complete sentences in grammatical English.

1) [30 points] Compute the lifetime of classical matter: at time  $t = 0$  an electron orbits a proton in a classical circular orbit of radius  $R$ . Compute how long it takes for the electron to spiral (non-relativistically) to  $R = 0$ , assuming that it always remains in a circular orbit as it radiates.

2) [30 points] The Greisen bound: there is thought to be an upper limit to the energy of galactic cosmic rays (protons) due to the inelastic scattering of protons off the 3 K cosmic microwave background,  $p + \gamma = \Delta$ . What is this bound? Use  $m_p = 1$  GeV,  $m_\Delta = 1.2$  GeV, and recall that room temperature (300 K) corresponds to a mean photon energy of 1/40 eV. As yet another hint, the bound is an energy much greater than the proton mass.

3) [30 points] Light of frequency  $\omega$  scatters from a sphere of radius  $R$  made of paramagnetic material, which is parametrized by a local fluctuation in the magnetic susceptibility  $\delta\mu$  which is a constant inside the sphere. (This is CGS language, in MKS,  $\epsilon = \epsilon_0$  and it is  $\mu/\mu_0$  which experiences a fluctuation away from unity.) Compute the (polarization summed and averaged) differential cross section in Born approximation for light scattering from the impurity.

4) [30 points] A particle moves relativistically along the  $x$  axis under a force due to a constant electric field  $\vec{F} = \hat{x}eE_0$ . It is at rest at  $x = 0$  at  $t = 0$ . (a) [20 points] What is  $x(t)$ ? Check your answer at very short and very long times. (b) [10 points] What is the total (instantaneous) electromagnetic power radiated at any time?

5) [30 points] A positron of initial velocity  $v/c = \beta_0$  collides with a nucleus of charge  $Ze$  and scatters directly backward along its initial direction of motion. Find a formula for  $d^2N/(d\omega d\Omega)$ , the number of photons emitted per solid angle  $d\Omega$  and frequency interval  $d\omega$ . Your answer should be valid when  $\beta_0$  is close to unity and at low  $\omega$ .

$$1) \quad P = -\frac{dE}{dt} = \frac{2}{3} \frac{e^2 a^2}{c^3}$$

1.1

For circular motion,  $E = \frac{1}{2} m v^2 - \frac{e^2}{r}$

$$\frac{m v^2}{r} = \frac{e^2}{r^2} \Rightarrow m v^2 = \frac{e^2}{r}$$

$$\rightarrow E = -\frac{1}{2} \frac{e^2}{r}$$

$$a = \frac{v^2}{r} = \frac{e^2}{m r^2} \quad \text{so} \quad P = \frac{2}{3} \frac{e^2}{c^3} \left( \frac{e^2}{m r^2} \right)^2$$

$$\frac{dE}{dt} = \frac{d}{dt} \left( -\frac{e^2}{2r} \right) = P \quad \text{so}$$

$$\frac{e^2}{2r^2} \frac{dr}{dt} = \frac{2}{3} \frac{e^2}{c^3} \left( \frac{e^2}{m r^2} \right)^2$$

$$\frac{dr}{dt} = \frac{4}{3} \frac{e^4}{m^2 c^3} \frac{1}{r^2}$$

$$\int_0^R r^2 dr = \frac{4}{3} \frac{e^4}{m^2 c^3} \int_0^T dt$$

$$\frac{1}{3} R^3 = \frac{4}{3} \frac{e^4}{m^2 c^3} T$$

$$\therefore T = \frac{1}{4} \frac{m^2 c^3}{e^4} R^3$$

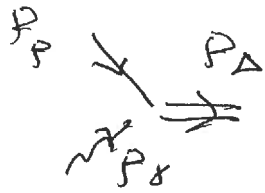
units:  $\left( \frac{\hbar c}{e^2} \right)^2 \frac{1}{(\hbar c)^2} m^2 c^4 \frac{R^3}{c} = \frac{\text{energy}^2}{(\text{energy} \times \text{length})^2} \frac{\text{length}^3}{\text{length} \times \text{time}}$

$$T = \frac{(m c^2)^2}{c} \left[ \frac{R}{e^2} \right]^2 \cdot R \sim \frac{(5 \times 10^5 \text{ eV})^2}{3 \times 10^{18} \text{ A/sec}} \frac{1}{13.6 \text{ eV}}$$

$$T \sim 10^{-8} \text{ sec.}$$

Greisen bound:  $p + \gamma = \Delta$  means

$$P_{\Delta} = P_p + P_{\gamma} \text{ as 4-vector relation}$$



$$\text{Square } \vec{\Delta}: P_{\Delta}^2 = m_{\Delta}^2 = (P_p + P_{\gamma})^2$$

$$m_{\Delta}^2 = m_p^2 + 2P_p \cdot P_{\gamma} \quad (P_{\gamma}^2 = 0)$$

$$\text{or } P_p \cdot P_{\gamma} = \frac{m_{\Delta}^2 - m_p^2}{2}$$

Now evaluate head-on collision in Earth's frame

$$P_{\gamma} = [\omega, 0, 0, -\omega] \quad \text{massless photon}$$

$$P_p = [E, 0, 0, E] \quad \text{very high energy, treat as massless particle}$$

$$P_{\gamma} \cdot P_p = 2\omega E \quad \text{or}$$

$$E = \frac{m_{\Delta}^2 - m_p^2}{4\omega}$$

Numbers: Room temp is 300 K, black body is 3 K,  $300 \text{ K} = \frac{1}{40} \text{ eV}$

$$\omega = \frac{1}{40} \text{ eV} \times \frac{1}{100} = \frac{1}{4000} \text{ eV}$$

$$m_{\Delta}^2 = 1.2 \text{ GeV}^2, \quad m_p^2 = 1 \text{ GeV}^2, \quad \text{GeV} = 10^9 \text{ eV}$$

$$m_{\Delta}^2 - m_p^2 = 0.44 \text{ GeV}^2 = 0.44 \times 10^{18} \text{ eV}^2$$

$$E = \frac{0.44 \times 10^{18} \text{ eV}^2}{\frac{1}{4} \times 10^{-3} \text{ eV}} \sim 10^{21} \text{ eV} !$$

Horrible un-necessary algebra

$$P_P \cdot P_\Delta = \omega \left[ \sqrt{E^2 - m^2} + E \right]$$

↑  
momentum

$$E + \sqrt{E^2 - m^2} = \frac{m_\Delta^2 - m_P^2}{2\omega}$$

$$\sqrt{E^2 - m^2} = \frac{m_\Delta^2 - m_P^2}{2\omega} - E$$

square:

$$E^2 - m^2 = \left[ \frac{m_\Delta^2 - m_P^2}{2\omega} \right]^2 - E \left( \frac{m_\Delta^2 - m_P^2}{2\omega} \right) + E^2$$

$$E = \frac{\left( \frac{m_\Delta^2 - m_P^2}{2\omega} \right)^2 - m^2}{\frac{m_\Delta^2 - m_P^2}{2\omega}}$$

The first term is huge because  $\omega \gg m_P$  vs  $m_P$ . ~~The first term~~ Neglect  $m_P$  to get

$$E \sim \frac{m_\Delta^2 - m_P^2}{2\omega} + \dots$$

the answer on p. 2.1.

3) This is a Born approximation problem, nearly done in the book. Jackson (10.3) says

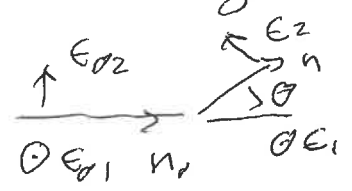
$$\frac{d\sigma}{d\Omega}(\vec{e}_0 \rightarrow \vec{e}) = |f(\theta)|^2$$

$$f(\theta) = \frac{k^2}{4\pi} (\hat{n} \times \hat{e}^*) \cdot (\hat{n}_0 \times \vec{e}_0) g(\beta)$$

$$g(\beta) = \int d^3x e^{i\vec{\beta} \cdot \vec{x}} \rho(x)$$

$$\frac{d\sigma}{d\Omega} = \frac{k^4}{16\pi^2} |g(\beta)|^2 \left| (\hat{n} \times \hat{e}^*) - (\hat{n} \times \vec{e}_0) \right|^2$$

We average initial polarizations, sum final ones



Note  $\hat{n} \times \hat{e}_2 = \hat{e}_1$  for each polarization  
 $\hat{n} \times \hat{e}_1 = -\hat{e}_2$

so the polarization part of the answer is like the  $\vec{e}_0 \cdot \vec{e}^*$  of electric scattering

$\vec{e}_0 \backslash \vec{e}$	1	2
1	$\cos^2 \theta$	0
2	0	1

$$\frac{d\bar{\sigma}}{d\Omega} = \frac{1}{2} \sum_{\vec{e}_0, \vec{e}} \frac{d\sigma}{d\Omega}$$

$$\frac{d\bar{\sigma}}{d\Omega} = \frac{k^4}{16\pi^2} \left[ \frac{1 + \cos^2 \theta}{2} \right] |g(\beta)|^2$$

$$g(\vec{q}) = \int d^3x e^{i\vec{q}\cdot\vec{x}} \cdot \delta_{\mu\nu}(v)$$

$$= \delta_{\mu\nu} \int_0^R r^2 dr \int_0^{2\pi} d\varphi \int_{-1}^1 d\cos\theta e^{i\vec{q}\cdot\vec{r} \cos\theta}$$

(Picking the z-axis to be along  $\vec{q}$ )

$$g = \delta_{\mu\nu} \cdot 2\pi \int_0^R r^2 dr \left[ \frac{e^{i\vec{q}\cdot\vec{r}} - e^{-i\vec{q}\cdot\vec{r}}}{i\vec{q}\cdot\vec{r}} \right]$$

$$= 4\pi \frac{\delta_{\mu\nu}}{\vec{q}} \int_0^R r dr \sin \vec{q}\cdot\vec{r}$$

$$= \frac{4\pi}{\vec{q}^3} \int_0^{\vec{q}\cdot\vec{R}} x dx \sin x = \frac{4\pi}{\vec{q}^3} \left[ \sin \vec{q}\cdot\vec{R} - \vec{q}\cdot\vec{R} \cos \vec{q}\cdot\vec{R} \right]$$

$$\frac{d\vec{r}}{d\Omega} = R^4 \left[ \frac{1 + \cos^2\theta}{2} \right] \left( \frac{\sin \vec{q}\cdot\vec{R} - \vec{q}\cdot\vec{R} \cos \vec{q}\cdot\vec{R}}{\vec{q}^3} \right)^2$$

4) Particle begins at rest in constant E-field

$$\frac{d}{dt} \gamma m v = e E \Rightarrow \gamma v = \frac{e E t}{m} \quad \frac{v}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$x(t) = \int v(t) dt.$$

$$\frac{v(t)/c}{\sqrt{1 - \frac{v(t)^2}{c^2}}} = \frac{e E t}{m c} \Rightarrow \frac{\beta^2}{1 - \beta^2} = \left( \frac{e E t}{m c} \right)^2$$

$$\beta^2 = (1 - \beta^2) \left[ \frac{e E t}{m c} \right]^2$$

$$\left[ 1 + \left( \frac{e E t}{m c} \right)^2 \right] \beta^2 = \left( \frac{e E t}{m c} \right)^2$$

$$\Rightarrow \beta(t) = \frac{v(t)}{c} = \frac{e E t / m c}{\left[ 1 + \left( \frac{e E t}{m c} \right)^2 \right]^{1/2}}$$

$$x(t) = \frac{e E}{m} \int_0^t \frac{t' dt'}{\left[ 1 + \left( \frac{e E t'}{m c} \right)^2 \right]^{1/2}} \quad \begin{matrix} u = 1 + \left( \frac{e E t'}{m c} \right)^2 \\ du = 2 t' \left( \frac{e E}{m c} \right)^2 \end{matrix}$$

$$x(t) = \frac{e E}{m} \left( \frac{m c}{e E} \right)^2 \cdot \frac{1}{2} \int \frac{du}{\sqrt{u}} = \frac{m c^2}{e E} \left[ 1 + \left( \frac{e E t'}{m c} \right)^2 \right]^{1/2} \Big|_0^t$$

$$x(t) = \frac{m c^2}{e E} \left[ \left( 1 + \left( \frac{e E t}{m c} \right)^2 \right)^{1/2} - 1 \right]$$

At small t,  $x(t) = \frac{m c^2}{e E} \left[ 1 + \frac{1}{2} \left( \frac{e E t}{m c} \right)^2 - 1 \right]$   
 $= \frac{e E}{m} \frac{1}{2} t^2$

At big t  $x(t) \sim \frac{m c^2}{e E} \left( \frac{e E t}{m c} \right) \sim c t!$

Power radiated?

4.2

$$P = -\frac{2}{3} \frac{e^2}{m^2 c^3} \frac{dP_{\mu}}{d\tau} \frac{dP^{\mu}}{d\tau}$$
$$= \frac{2}{3} \frac{e^2}{m^2 c^3} \gamma^2 \left[ \left( \frac{d\vec{p}}{dt} \right)^2 - \frac{1}{c^2} \left( \frac{dE}{dt} \right)^2 \right]$$

$$\frac{d\vec{p}}{dt} = e \mathbf{E} \text{ from electric field } E$$

energy  $E^2 = p^2 c^2 + m^2 c^4 \Rightarrow$

$$\frac{E}{c} \frac{dE}{dt} = \frac{c^2 p}{c} \frac{dp}{dt}$$

$$\frac{1}{c} \frac{dE}{dt} = c \frac{p}{E} \frac{dp}{dt} \quad ; \quad \frac{c p}{E} = \beta = \frac{v}{c} \Rightarrow$$

$$\frac{1}{c} \frac{dE}{dt} = \frac{v}{c} \frac{dp}{dt}$$


$$P = \frac{2}{3} \frac{e^2}{m^2 c^3} \left( \frac{dp}{dt} \right)^2 \cdot \gamma^2 \left( 1 - \frac{v^2}{c^2} \right)$$
$$= \frac{2}{3} \frac{e^2}{m^2 c^3} \left( \frac{dp}{dt} \right)^2 \quad (\text{This is 14.27!})$$
$$= \frac{2}{3} \frac{e^2}{m^2 c^3} (eE)^2$$



5) Positron collides with nucleus. In the extreme relativistic limit and at low frequencies and small angles,

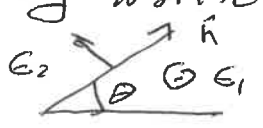
$$\frac{d^2 I}{d\Omega d\omega} = \frac{e^2}{4\pi^2 c} \left| \hat{\epsilon}^* \cdot \left( \frac{\vec{\beta}_+}{1 - \vec{\beta}_+ \cdot \hat{n}} - \frac{\vec{\beta}_-}{1 - \vec{\beta}_- \cdot \hat{n}} \right) \right|^2$$

$\vec{\beta}_- = \vec{\beta}_0$ ,  $\vec{\beta}_+ = -\vec{\beta}_0$ : let them lie on the Z axis



$$\left( \right) = - \left( \frac{1}{1 + \beta_0 \cos \theta} + \frac{1}{1 - \beta_0 \cos \theta} \right) \vec{\beta}_0 = - \frac{2\vec{\beta}_0}{1 - \beta_0^2 \cos^2 \theta}$$

To make life easy while summing polarizations, pick them as  $\hat{\epsilon}_1, \hat{\epsilon}_2$  so  $\hat{\epsilon}_1 \cdot \vec{\beta}_0 = 0$ ,  $\hat{\epsilon}_2 \cdot \vec{\beta}_0 = \beta_0 \sin \theta$



$$\frac{d^2 I}{d\Omega d\omega} = \frac{e^2}{4\pi^2 c} \cdot \frac{4\beta_0^2 \sin^2 \theta}{(1 - \beta_0^2 \cos^2 \theta)^2}$$

$I(\omega) = \hbar \omega N(\omega)$  so

$$\frac{d^2 N}{d\Omega d\omega} = \frac{e^2}{\hbar c} \frac{1}{4\pi^2} \frac{1}{\omega} \cdot \frac{4\beta_0^2 \sin^2 \theta}{(1 - \beta_0^2 \cos^2 \theta)^2}$$

Good thing:  $\beta_0^2 = 1 - \frac{1}{\gamma^2}$ ,  $\sin \theta \sim \theta$ ,  $\cos^2 \theta = 1 - \theta^2$

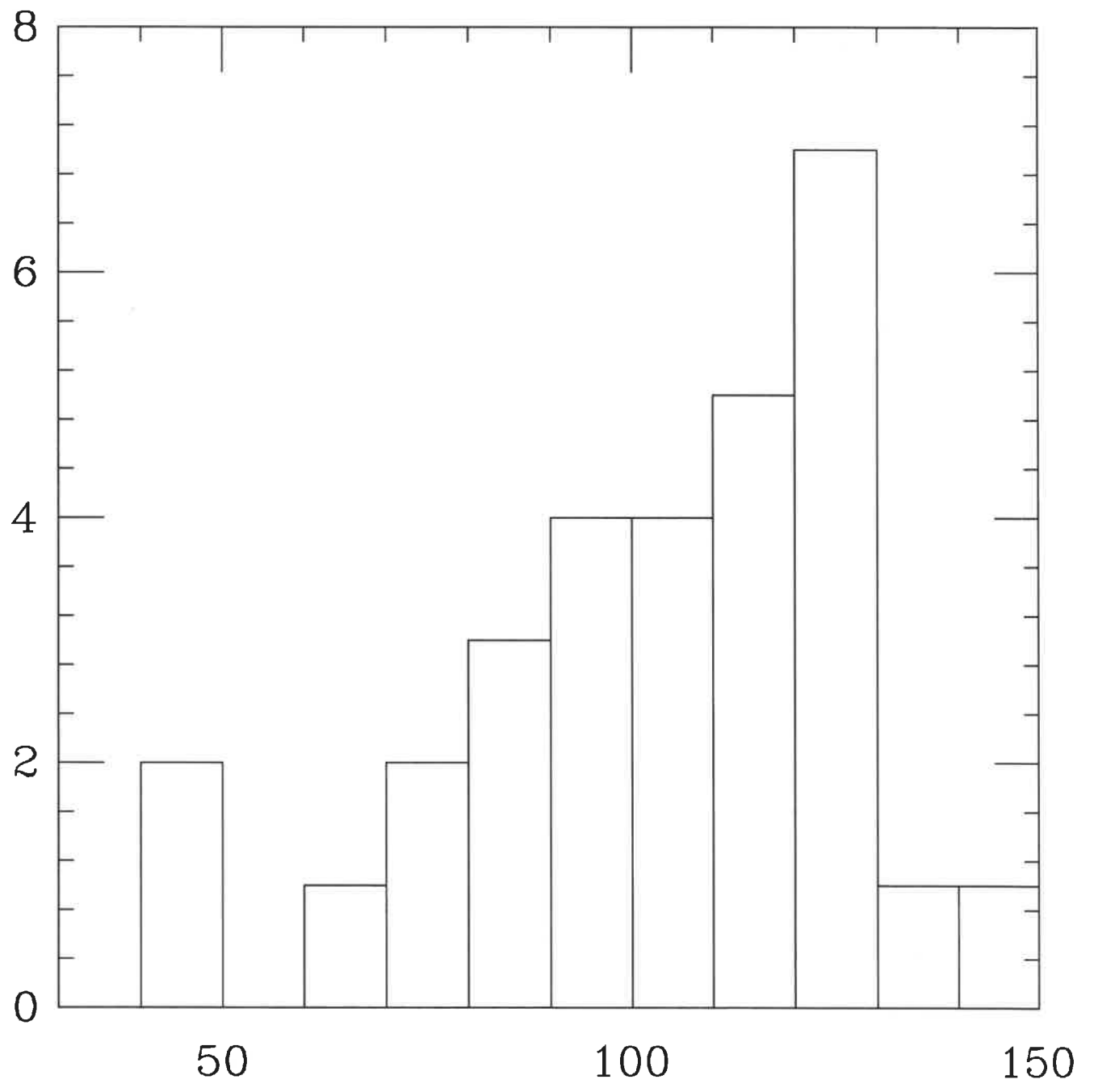
$$1 - \beta_0^2 \cos^2 \theta = 1 - \left(1 - \frac{1}{\gamma^2}\right) (1 - \theta^2) = \frac{1}{\gamma^2} + \theta^2 = \frac{1}{\gamma^2} (1 + \gamma^2 \theta^2)$$

$$\frac{d^2 N}{d\Omega d\hbar\omega} = \frac{e^2}{\hbar c} \frac{1}{\pi^2} \frac{1}{\hbar\omega} \cdot \gamma^2 \left[ \frac{(\gamma\theta)^2}{(1 + \gamma^2 \theta^2)^2} + (\theta \rightarrow \pi - \theta \text{ for backwards terms}) \right]$$

7320 Spring 2024

Final  $\mu = 101$

$\sigma = 25$



7320 overall  
Spring 2024

ave = 382/500  
 $\sigma = 37$

