

Set 1– due 8 September

“You won’t become a good pianist by listening to good concerts.” – J. Wess

For this week, we have some short problems, while we build tools: First some (rather formal) Green’s theorem manipulations...

1) [5 points] Jackson 1.10: note that this is only true if there is no charge inside the sphere!

2) [5 points] Jackson 1.12: this is needed for

3) [10 points] Jackson 1.13

4) [20 points] A sphere of radius R has a uniform charge density ρ and total charge Q . Find the electrostatic potential energy of the sphere, that is, the energy required from some non-electric interaction to hold the sphere together. On dimensional grounds, the answer will be

$$W = C \frac{1}{4\pi\epsilon_0} \frac{Q^2}{R} \quad (1)$$

and the interesting quantity is the constant C . Do this problem “like an undergraduate,” that is, use Gauss’s law to find \vec{E} , then integrate it to find $\Phi(r)$. Do the problem two ways: compute (a) [10 points]

$$W = \frac{1}{2} \int d^3x \rho(x) \Phi(x) \quad (2)$$

(b) [10 points]

$$W = \frac{\epsilon_0}{2} \int d^3x |\vec{E}(x)|^2. \quad (3)$$

We’ll visit the third alternative

$$W = \frac{1}{2} \int \rho(r) \rho(r') \frac{1}{4\pi\epsilon_0} \frac{1}{|\vec{r} - \vec{r}'|} d^3r d^3r' \quad (4)$$

later in the course—this looks like a lot but it is actually the easiest way to do the problem.

Jackson 1.10 . From Jackson eq -136 we have

1-1

$$\Phi(x) = \frac{1}{4\pi\epsilon_0} \int_V d^3x' \frac{\rho(x')}{|x-x'|} + \frac{1}{4\pi} \int_S dA' \left[\frac{1}{R} \frac{\partial \Phi}{\partial n'} - \Phi \frac{\partial}{\partial n'} \frac{1}{R} \right]$$

If $\rho(x')=0$ the first term vanishes. Let S be a sphere of radius R . Let x be a point on the interior - in fact let it be in the center and choose conveniently $\rho(x)=0$. Then \hat{n}' and x' are parallel, $r'=|\vec{x}'|=R$

$$\frac{\partial \Phi}{\partial n'} = -\vec{E} \cdot \hat{n}' \rightarrow \frac{\partial}{\partial r'} \frac{1}{r'} \Big|_{r'=R} = -\frac{1}{R^2}$$

using $\frac{\partial}{\partial n'} = \hat{n}' \cdot \nabla = \frac{\partial}{\partial r'}$. Put everything together

$$\Phi(0) = -\frac{1}{4\pi R} \int_S dA' \vec{E} \cdot \hat{n}' + \frac{1}{4\pi R^2} \int_S dA' \Phi(x')$$

Gauss' law says $\int dA' \vec{E} \cdot \hat{n}' = \frac{Q_{\text{enclosed}}}{\epsilon_0} = 0$

$$\text{so } \Phi(0) = \frac{1}{4\pi R^2} \int_S dA' \Phi(x')$$

The potential at the center is the average value of Φ on the surface of the sphere.

2) 1-12 Call the potential due to $\epsilon \neq \sigma$ $\underline{\Phi}$
 call the potential due to $\epsilon' \neq \sigma'$ $\underline{\Phi}'$.
 We are asked to show that

$$\int_V \epsilon \underline{\Phi}' d^3x + \int_S \sigma \underline{\Phi}' dA = \int_{V'} \epsilon' \underline{\Phi} d^3x + \int_S' \sigma' \underline{\Phi} dA$$

Start with Jackson eq. 1.35, Green's 2nd identity,

$$\text{set } \phi = \underline{\Phi} \quad \psi = \underline{\Phi}'$$

$$\epsilon = -\epsilon_0 \nabla^2 \underline{\Phi} \quad \epsilon' = -\epsilon_0 \nabla^2 \underline{\Phi}'.$$

Also recall that outside a conductor

$$\frac{\partial \underline{\Phi}}{\partial n} = -E_n = -\frac{\sigma}{\epsilon_0} - \underline{\text{B}}_{\text{ext}} \quad (\text{in this form only})$$

\hat{n} is an outward normal  while in

1.35 \hat{n} is a normal into the integration region.

$$1.35: \int_V (\phi \nabla^2 \psi - \psi \nabla^2 \phi) d^3x = \int_S dA \left[\phi \frac{\partial \psi}{\partial n} - \psi \frac{\partial \phi}{\partial n} \right]$$

$$\text{or } -\frac{1}{\epsilon_0} \int_V (\underline{\Phi} \epsilon' - \underline{\Phi}' \epsilon) d^3x = -\frac{1}{\epsilon_0} \int_S dA [-\underline{\Phi} \sigma' + \underline{\Phi}' \sigma]$$

depth of normal

Just regroup

$$\frac{1}{\epsilon_0} \int_V \underline{\Phi} \epsilon' d^3x + \int_S \underline{\Phi} \sigma' dA = \int_{V'} \underline{\Phi}' \epsilon d^3x + \int_S' \underline{\Phi}' \sigma dA$$

Now we apply it---

unprimed system



$$\bullet x = x_0$$



$$\sigma(x) = q \delta^3(x - x_0)$$

$$\vec{x}_0 = (x_0, 0, 0)$$

$\Phi = 0$ on both surfaces

$\sigma = \sigma(d)$ on the top plate
also σ on (own plate)

Both systems are infinite in the transverse directions -
we need this to be true to solve the primed system.

Reciprocity formula says

$$q \int_V \delta(x - x_0) \delta(y) \delta(z) \left[\frac{\sigma_0 x}{\epsilon_0} \right] d^3x + \underbrace{\left(\Phi' \sigma + \int_{S_-} \sigma_0 \cdot \sigma \cdot dA + \int_{S_+} \cancel{\sigma(d)} \sigma(d) \frac{\sigma_0 d}{\epsilon_0} dA \right)}_{\Phi' \sigma}$$

$$= \int_V \sigma \cdot \Phi d^3x + \int_S \sigma' \cdot \sigma dA = 0$$

$\Sigma \Phi' \sigma$ $\Sigma \sigma' \sigma$

$$\text{or } q \frac{\sigma_0 x_0}{\epsilon_0} + \frac{\sigma_0 d}{\epsilon_0} \int_{S_+} \sigma(d) dA = 0$$

$$-\frac{x_0}{d} q = \int_{S_+} \sigma(d) dA = \text{induced charge on } S_+ - \text{We'll derive this less mechanically later in the term.}$$

$$4). W = \frac{1}{2} \int \epsilon_0 E^2 d^3x \quad \text{or} \quad \frac{\epsilon_0}{2} \int E^2 d^3a \quad L4-1$$

Gauss: for $r < R$ $4\pi r^2 E = \frac{1}{\epsilon_0} \left(\frac{r^3}{R^3} Q \right)$

$$\left(\oint E \cdot d\vec{l} = \frac{Q}{\epsilon_0 R^3} \int_0^r 4\pi r'^2 dr' \right)$$

$$\text{so } E_r = \frac{1}{4\pi\epsilon_0} \frac{Q}{R^2} \quad \text{inside}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \quad \text{outside}$$

$$\Phi = \frac{Q}{4\pi\epsilon_0} \frac{1}{r} \quad \text{outside as } \Phi \rightarrow 0 \text{ as } r \rightarrow \infty.$$

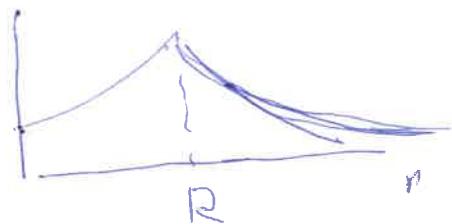
Then inside $\Phi(r) - \Phi(0) = - \int_0^r E \cdot dr = - \frac{Q}{4\pi\epsilon_0} \frac{1}{R^2} \int_0^r r dr$

$$= - \frac{Q}{4\pi\epsilon_0} \frac{1}{R^3} \cdot \frac{1}{2} r^2$$

$$\Phi(R) = \Phi(0) - \frac{Q}{4\pi\epsilon_0} \frac{1}{2} \frac{R^2}{R^3} = \frac{1}{4\pi\epsilon_0} \frac{Q}{R}$$

$$\text{so } \Phi(0) = \frac{1}{4\pi\epsilon_0} \left[\frac{Q}{R} + \frac{Q}{2R} \right] = \frac{3}{2} \frac{Q}{4\pi\epsilon_0} \frac{1}{R}$$

$$\text{or } \Phi(r) = \frac{Q}{4\pi\epsilon_0} \left[-\frac{1}{2} \frac{r^2}{R^3} + \frac{3}{2} \frac{1}{R} \right] \quad \text{if } r < R$$



Part (a)

$$\begin{aligned}
 W &= \frac{1}{2} \int_0^R 4\pi r^2 dr \epsilon(r) \Phi(r) \\
 &= \frac{1}{2} \left[\frac{\Phi}{\frac{4\pi R^3}{3}} \right] 4\pi \int_0^R r^2 dr \cdot \frac{\Phi}{4\pi \epsilon_0} \left(-\frac{1}{2} \frac{r^2}{R^3} + \frac{3}{2} \frac{1}{R} \right) \\
 &= \frac{\Phi^2}{4\pi \epsilon_0} \cdot \frac{3}{2} \cdot \frac{1}{R^3} \cdot \int_0^R dr \left(-\frac{1}{2} \frac{r^4}{R^3} + \frac{3}{2} \frac{r^2}{R} \right) \\
 &= \frac{\Phi^2}{4\pi \epsilon_0} \cdot \frac{3}{2} \frac{1}{R^3} \left\{ -\frac{1}{2} \cdot \frac{1}{5} \frac{R^5}{R^3} + \frac{3}{2} \cdot \frac{1}{3} \frac{R^3}{R} \right\} \\
 &= \frac{\Phi^2}{4\pi \epsilon_0} \frac{1}{R} \cdot \frac{1}{9} \\
 \frac{1}{9} &= \frac{3}{2} \left\{ -\frac{1}{10} + \frac{1}{2} \right\} = \frac{3}{2} \cdot \frac{4}{10} = \underline{\underline{\frac{3}{5}}}
 \end{aligned}$$

4.3

$$b) W = \frac{\epsilon_0}{2} \int d^3x E^2(x)$$

$$= \frac{\epsilon_0}{2} \int_{r < R} 4\pi r^2 dr E_m^2 + \frac{\epsilon_0}{2} \int_{r > R} 4\pi r^2 dr E_{out}^2$$

$$= \frac{\epsilon_0}{2} \cdot \left(\frac{1}{4\pi\epsilon_0} \right)^2 4\pi - Q^2 \times$$

$$\left\{ \int_0^R r^2 dr \left(\frac{r}{R^3} \right)^2 + \int_R^\infty r^2 dr \frac{1}{r^4} \right\}$$

$$= \frac{Q^2}{4\pi\epsilon_0} \cdot \frac{1}{2} \cdot \left[\frac{1}{5} \frac{R^5}{R^6} + \frac{1}{R} \right]$$

$$= \frac{Q^2}{4\pi\epsilon_0} \cdot \frac{1}{R} \cdot \frac{1}{5}$$

$$\frac{1}{5} = \frac{1}{2} \left(\frac{1}{5} + 1 \right) = \frac{1}{2} \cdot \frac{6}{5} = \frac{3}{5} \text{ again!}$$

